

Study of interference effects on the photoluminescence of AlGa_N/Ga_N quantum wells

M. Ramírez-López^{1,3}, Y. L. Casallas-Moreno¹, M. Pérez-Caro¹, A. Escobosa-Echevarria, S. Gallardo-Hernández^{*,1,2}, J. Huerta-Ruelas⁴, and M. Lopez-Lopez¹

¹ Physics Department, Centro de Investigación y Estudios Avanzados del IPN, Apartado Postal 14-740, México D.F., México 07360

² Electrical Engineering Department, Centro de Investigación y Estudios Avanzados del IPN, Apartado Postal 14-740, México D.F., México 07360

³ Unidad Profesional Interdisciplinaria en Ingeniería y Tecnologías Avanzadas del IPN, Av. Instituto Politécnico Nacional 2580, México 07340

⁴ Centro de Investigación en Ciencia Aplicada y Tecnología Avanzada del IPN, Cerro Blanco 141, Querétaro, México 76090

Received 14 July 2014, revised 28 October 2015, accepted 20 January 2015

Published online 19 March 2015

Keywords Ga_N, PAMBE, modulated photoluminescence

* Corresponding author: e-mail sgallardo@fis.cinvestav.mx, Phone: +52 55 5747 3828

In this work we compared the optical properties of Al-GaN(50 nm)/Ga_N(5 nm) quantum wells (QWs) grown by plasma-assisted molecular beam epitaxy (PAMBE) technique, on Si(111) substrates and sapphire substrates. Optical properties were acquired by photoluminescence (PL) and reflectance (R) spectroscopies. Reflectance spectra shows interference oscillations of the reflected beam at Air/AlGa_N and the beam reflected at bottom interface with Si substrate. Such oscillations vanish at the band gap of Ga_N buffer layer (3.4 eV). The PL spectra of samples grown on Silicon substrate shows a photolumi-

nescence modulation effect, attributed to interference of light emitted from the QWs that is reflected at different heterostructure interfaces. On sapphire substrate no modulation effect is present. This could be explained due to interfacial roughness and the smaller refractive index of sapphire in comparison to silicon. PL shows a strong emission around 3 eV, which is in agreement with recombination energy determined by self-consistent calculations, which consider a 4 MV/cm built-in electric field and low carrier densities.

© 2015 WILEY-VCH Verlag GmbH & Co. KGaA, Weinheim

1 Introduction Nitride semiconductor materials have enabled the design of nanostructures that can operate optically from the infrared to the ultraviolet of the electromagnetic spectrum, which are highly suitable for making devices like electroluminescent diodes, UV sensors and high efficiency solar cells of second and third generation [1]. Some of these devices are grown with microcavities to enhance optical characteristics like the output power [2]. Most common microcavities are based on super-lattices, however few years ago, some researchers propose simple microcavities designed to match with the wavelength of the incident light [3]. In this work we present the results of Fabry-Perot effects due to an inherent microcavity without the contribution of a top Bragg reflector. Our structure could be potentially employed for improvement of ultra-

violet Light Emission Diodes with a Ga_N active layer using AlGa_N barriers.

2 Experimental The structures were grown on a Riber C21 molecular beam epitaxy (MBE) system. Active Nitrogen (N) flux was provided by an rf-plasma source operated at 150 W with a N₂ flow of 0.25 sccm. We employed two types of substrates: silicon and sapphire. Ratio of III/N fluxes was higher than one to ensure a metallic rich growth condition that leads to a smoother surface, a decrease in defect density and higher electron mobility [4]. In the case of silicon substrates, they were previously chemically cleaned by employing the Ishizaka and Shiraki method [5]. Surface oxides on substrates were thermally desorbed at 900 °C. Next, substrate temperature was decreased to 750 °C, at this temperature a 7x7 surface recon-

struction was observed by Reflection High-Energy Electron Diffraction (RHEED) indicating a successful oxide desorption. Then a metallization of the Si surface with an aluminum layer (~5 monolayers) was performed at 850 °C during the onset of nitrogen plasma source. This was done to avoid nitrogen incorporation and the formation of SiN which involves amorphization of Si surface. After this process an AlN layer was grown at 850 °C with a nominal thickness of 27 nm. After that, a 5 period superlattice of AlN/GaN was grown to get a smoother surface, followed by a GaN buffer layer with a thickness of 600 nm employing a substrate temperature of 800 °C. Finally, a series of 5 quantum wells were grown with a GaN 5nm-thick active layer and 50 nm-thick AlGaIn barriers with an aluminum composition of 30%. A growth stop of 80 s was applied after each of the five GaN active layers. For the growth on sapphire substrates, we performed a surface nitridation process at 900 °C for 20 min. Then substrate temperature was decreased to 800 °C to grow an AlN buffer layer, followed by the same growth procedure as for Si substrates.

3 Results Figure 1 shows the intensity behavior of RHEED specular spot during QWs growth. We can observe that the RHEED intensity decreases just after the Ga shutter is opened, which indicates an increase in surface roughness. When Ga shutter is closed, the RHEED intensity recovered during the 80 s growth stop, which suggests a smoothing of the surface. At the end of the growth we observed a (2x2) RHEED pattern (inset in Fig. 1), which confirms an atomically flat surface with Ga polarity [6].

The samples were characterized structurally by high resolution X-ray diffraction (HRXRD) on a Panalytical MRD Xpert diffractometer. We employed K_{α} radiation of Cu and spectra were obtained with a two bounced monochromator and a three bounced analyzer in a triple axis arrangement. Diffraction patterns were obtained to verify the formation of wurtzite structure for GaN and AlGaIn layers. Also we obtained rocking curves for the diffraction of {0002} planes of GaN, AlGaIn and AlN. From the full width at half maximum (FWHM) of the rocking curves (Table 1) we found that crystalline quality is slightly better for growth over silicon substrates.

Table 1 FWHM of HRXRD rocking curves of {0002} planes of GaN, AlGaIn and AlN on Si and Sapphire substrates.

Substrate	Si	Sapphire
GaN	0.47	0.49
AlN	0.63	0.69
AlGaIn	0.59	0.51

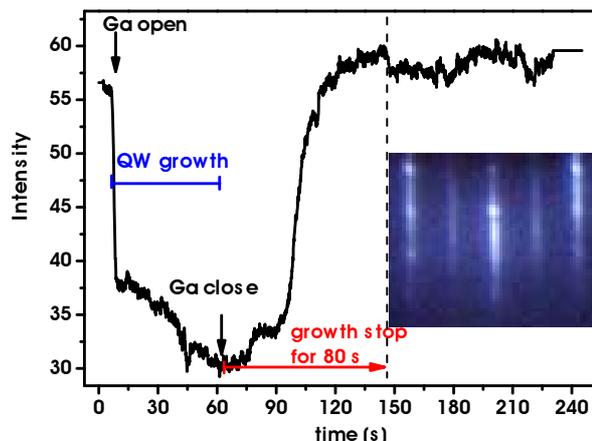


Figure 1 RHEED intensity during growth of the GaN active layer of quantum wells.

Photoluminescence (PL) signal was obtained by optical excitation of samples with a He-Cd laser and luminescence was analyzed employing a Horiba/Spex 1404 0.85m Double Spectrometer. Sample grown on Si(111) substrate presents a PL spectrum centered at 3 eV, which is modulated by oscillations that coincides with its reflectance spectrum (Fig. 2). For sapphire substrate, PL spectrum is centered at 3.1 eV without oscillations.

For both samples, we obtained a strong red shift in spectra due to combined effect of internal electric field and the carrier distribution. By using the mode spacing expression as in Ref. [2], we calculate a 700 nm thickness of the vertical cavity formed between the air/AlGaIn and the GaN/AlN interfaces, considering the peaks located at 2.9 and 3.1 eV, respectively.

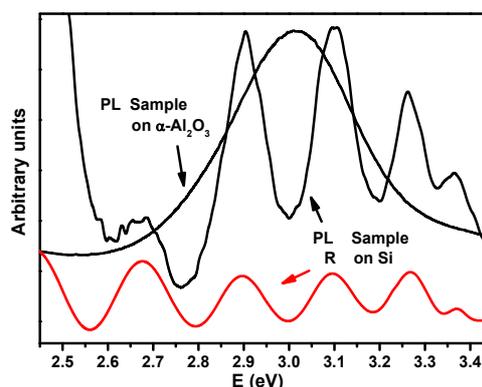


Figure 2 PL of AlGaIn/GaN QWs grown on Si (red) and sapphire (black). Reflectance of QWs grown on Si (blue).

For further investigation of the modulation effect on the photoluminescence spectrum, we measured PL and R spectra at different emission angles relative to sample normal and different incident angles as shown in Fig. 3. Maximum modulation peaks present a shift to higher energies as a function of incidence angle.

Equation (1) describes the energy position dependence of constructive interference as a function of the measurement angle. From this equation, when the angle θ_m at which the spectrum is measured increases, energy of constructive interference increase, because $\cos(\theta_m)$ decrease.

$$h\nu = \frac{(2m+1)hc}{d_{\text{buffer}}n(\lambda)\cos(\theta_m)}, \quad (1)$$

where d_{buffer} is the thickness of microcavity, n the refractive index and $(2m+1)$ are integers which define maximal reflectivity points [7].

In order to determine the theoretical position of luminescence, we calculated a band diagram for a single quantum well obtained by self-consistent solution of Schrödinger and Poisson's equations by employing a variational method, which can be resumed in the following equations:

$$\sigma_{\text{pol}} = P_{\text{tot,layer1}} - P_{\text{tot,layer2}}, \quad (2)$$

$$\sigma_{\text{pol}} = (P_{\text{SP}} + P_{\text{PE}})_{\text{GaN}} - (P_{\text{SP}} + P_{\text{PE}})_{\text{AlGaIn}}, \quad (3)$$

$$h\nu = E_g + E_e + E_{hh} + \delta_{E_e} + \delta_{E_{hh}}, \quad (4)$$

$$\delta E_i = \frac{2maeE}{\hbar^2} z^{22}, \quad (5)$$

where E_g is the gap of the semiconductor QW, $E_{e(h)}$ is the electron (hole) energy without electric field in the QW and $\delta_{E_{e(h)}} = H_{e(h)} - E_{e(h)}$ is the induced energy shift for electrons (holes) which depends on the carrier density and the built-in electric field [8].

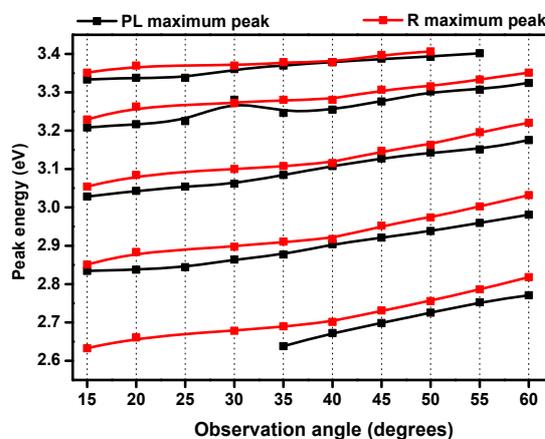


Figure 3 Angle dependence of PL (black squares) and reflectance (red circles) maximum for sample grown on Si.

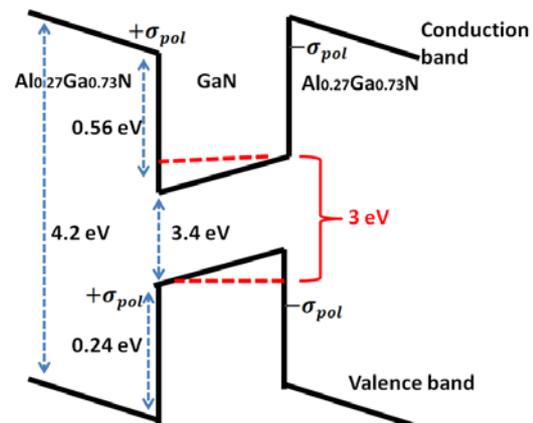


Figure 4 Band diagram of a single quantum well using variational method.

In the construction of the band diagram we consider a band gap of 3.4 eV for GaN and 4.4 eV for AlGaIn, and a low density of carriers regime ($\eta < 10^{10} \text{ cm}^{-2}$) at which, it is not possible to get screening of internal electrical fields. In addition to this, the AlGaIn barriers must growth under biaxial tension, due to a shorter lattice parameter than GaN. The barrier will show spontaneous polarization (P_{sp}) and piezoelectric (P_{pz}), while QW only presents spontaneous polarization. (P_{sp}). The results of such calculation are summarized in Fig. 4. From this model we obtained a recombination energy $E_{\text{PL}} = 3 \text{ eV}$, which correspond to the transition energy between the highest hole level and the lowest electron level. This occurs when the band offset is 0.7 eV for the conduction band and an internal piezoelectric field is 3.3 MW/cm.

4 Discussion Fabry-Perot interferences effects have been reported previously in InGaIn/GaN heterostructures [9]. In our case, we have proven to get a Fabry-Perot interference using a thinner structure based in a AlGaIn/GaN layers. We think that the GaN/AlIn superlattice at the very beginning of the whole structure could contribute to obtain a higher reflection of light generated at QWs during PL excitation. In order to theoretically verify this hypothesis, we calculated the reflectivity for such structure using an expression which do not depend on layer thickness but wavelength [10]. We found that reflectivity percentage of the rays originated at QWs for this superlattice structure is 60% at 325 nm wavelength. Another possible effect of the superlattice is to decrease the strain of GaN buffer layer, and also to smooth out the surface roughness of AlN layer. Here it is important to comment that we grew the same structure but without the GaN/AlIn superlattice and found no modulation effects on luminescence.

We neither found PL modulation effects on set of samples grown over silicon substrates with similar conditions, but employing lower substrate temperatures of 780 °C and 760 °C. The interference effects are present in samples, grown with the highest substrate temperature. It is clear

that with decrease of substrate temperature, we decrease the surface mobility of adatoms during the growth. For such reasons, we think that this effect is very sensitive to interface roughness

5 Conclusion We have found Fabry-Perot interference effects in GaN/AlGaIn QWs grown on Si substrates. The PL interference peaks have a shift to higher energies when the angle Θ_m at which the spectrum is measured increases, according to the interference condition. The photoluminescence modulation is due to the Fabry-Perot interference of the emitted light from the QWs when it is reflected at the AlN/GaN SL/Si(111) and AlGaIn/air interfaces. We consider such effect is not present in sample grown on sapphire due to higher roughness of AlN/sapphire interface, which is more dispersive, and also because the effective transmission coefficient for sapphire substrate is higher than silicon.

Acknowledgements This work was partially supported by CONACYT-SENER project No. 151076. We also like to thanks to A. Tavira for HRXRD rocking curve measurements.

References

- [1] J. Wu, *J. Appl. Phys.* **106**, 01110 (2009).
- [2] N. Nakada et al., *Appl. Phys. Lett.* **76**, 1804 (2000).
- [3] F. Sermond et al., *Appl. Phys. Lett.* **87**, 021102 (2005).
- [4] C. Adelman et al., *J. Appl. Phys.* **91**, 9638 (2002).
- [5] A. Ishizaka and Y. Sharaki, *J. Electrochem. Soc.* **133**, 666 (1986).
- [6] C. T. Foxon et al., *J. Cryst. Growth* **207**, 1-7 (1999).
- [7] E. Hecht and A. Zajac, *Optics* (Addison-Wesley, 1974), p. 306.
- [8] M. Ramírez-López et al., *Microelectron. J.* **39**, 447 (2008).
- [9] C. Hums et al., *J. Appl. Phys.* **101**, 033113 (2007).
- [10] C. J. R. Sheppard, *Pure Appl. Opt.* **4**, 665 (1995).